Glucose Metabolism and Oscillatory Behavior of Pancreatic Islets

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A variety of oscillations are observed in pancreatic islets. We establish a model, incorporating two oscillatory systems of different time scales: One is the well-known bursting model in pancreatic β -cells and the other is the glucose-insulin feedback model which considers direct and indirect feedback of secreted insulin. These two are coupled to interact with each other in the combined model, and two basic assumptions are made on the basis of biological observations: The conductance $g_{K(ATP)}$ for the ATP-dependent potassium current is a decreasing function of the glucose concentration whereas the insulin secretion rate is given by a function of the intracellular calcium concentration. Obtained via extensive numerical simulations are complex oscillations including clusters of bursts, slow and fast calcium oscillations, and so on. We also consider how the intracellular glucose concentration depends upon the extracellular glucose concentration, and examine the inhibitory effects of insulin.

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I. INTRODUCTION

Insulin secretion from pancreatic β cells is critical for glucose homeostasis in blood. Perturbations from the basal glucose concentration induce various oscillations with different periods in pancreatic β cells: fast oscillations with periods of several seconds [1], slow oscillations with periods of several minutes [2], and ultradian oscillations with periods of a few hours [3]. Each type of oscillations has been studied separately, within the respective model according to each time scale. For example, slow oscillations can be caused by a feedback mechanism having slow reaction time [2]. Independently of this, a microscopic model for the bursting mechanism of the membrane potential in a β cell was devised and investigated [1] while the bursting activity was also demonstrated in a simple two-dimensional map [4]. In particular, the microscopic bursting model, which was constructed in a similar manner to the Hodgkin-Huxley model for nerve excitation [5], was improved by the introduction of the potassium ion K⁺ conductance depending dynamically on the ATP concentration [6]. However, the correlations between the bursting mechanism and insulin secretion have not been addressed in an appropriate way and there still lacks a model describing the complex phenomena such as clusters of bursting action potentials [7], fast and slow Ca²⁺ oscillations [8], which arise from combined behaviors of different time scales. The purpose of this study is to propose a combined model, explaining both fast and slow oscillations of action potentials.

On the other hand, recent experimental investigations in vivo and in vitro about the oscillations of insulin secretion have revealed a more complex picture on the molecular basis. Stimulation of insulin secretion by glucose involves a rise in the cytoplasmic concentration of calcium ions (Ca²⁺) in β cells. This rise essentially re-

sults from the following sequence of events, as sketched in Fig. 1: Glucose transported by GLUT-1 and GLUT-2 transporters in a β cell raises the ratio of adenosine triphosphate (ATP) to adenosine diphosphate (ADP), which promotes closure of ATP-sensitive K⁺ channels. This generates membrane depolarization, urging voltagedependent calcium channels to open. The subsequent increase in free cytosolic Ca²⁺ then stimulates insulin secretion [9]. Note also that there are several experimental evidences for insulin to inhibit bursting of β cells [10, 11]: in particular, there is a report that insulin activates ATPsensitive K^+ channels in β cells [12]. Based on these experiments, we assume two possible inhibitory pathways of insulin: indirect effects on gating of K(ATP) channels by inhibiting glucose transport through GLUT-2 transporters by insulin and direct effects of insulin to activate the channel. Both negative feedback processes of insulin to activate the channels may result in oscillatory behavior. Oscillations of the membrane potential [13] drive oscillations of Ca²⁺ [14], leading in turn to oscillations of insulin secretion [15, 16]. In fact, the measured portalvein insulin secretion rate shows periodic behavior of insulin secretion with successive secretion and rest. The frequency of such pulsatile insulin secretion in vivo and that in vitro in an isolated perfused pancreas have been reported to vary from four to fifteen minutes and from six to ten minutes, respectively [17, 18].

These experimental results make it desirable to establish a model that can describe the whole feedback process on the microscopic level, with the aforementioned macroscopic features incorporated. As an attempt toward such a goal, we propose a model, connecting the macroscopic description of glucose regulation [2] and the microscopic mechanism of bursting behavior of β cells [1] via the prescription based on experimental observations: The conductance $g_{K(ATP)}$ for the K(ATP) current, i.e.,

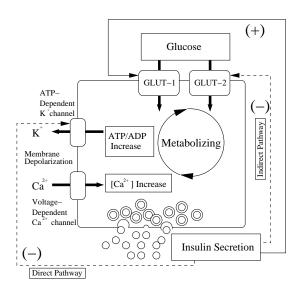


FIG. 1: Mechanism of insulin secretion. Activation and inhibition of GLUT-1 and GLUT-2 transporters by secreted insulin are represented by the solid (+) and dashed (-) arrows. Thick arrows describe physical transport of materials (glucose and ions).

the K⁺ current through the ATP-dependent K⁺ channel, is taken to be a decreasing function of the intracellular glucose concentration whereas it is assumed that the insulin secretion rate depends on the intracellular Ca²⁺ concentration. For the macroscopic description, we consider two possible insulin induced inhibitory pathways affecting the conductance $g_{K(ATP)}$. In addition, electrical coupling through gap junctions between β cells in an islet is assumed. Obtained via extensive numerical simulations are a variety of oscillatory behavior of the bursting electrical activity, calcium concentration, and insulin secretion, which are in good agreement with those observed experimentally.

There are four sections in this paper: Section II introduces the model system, described by coupled differential equations, together with the appropriate parameters. The coupled equations are integrated numerically and the results are presented in Sec. III. Finally, a detailed discussion, together with a brief summary, is given in Sec. IV.

II. MODEL SYSTEM

To describe the whole feedback process, we incorporate the bursting mechanism of β cells given by Sherman [1] and the description of glucose-insulin oscillations by Maki and Keizer [2, 19]. We begin with the dynamic equation for the time-dependent behavior of the membrane poten-

tial V_i of the *i*th β cell [1]:

$$\tau \frac{dV_i}{dt} = -I_{iCa} - I_{iK} - I_{iS} - I_{iK(ATP)} - \sum_{j}' g_c(V_i - V_j),$$
 (1)

which describes current balance with the relaxation time τ . The right-hand side of Eq. (1) includes contributions from various current channels, which are all given in units of voltage (i.e., with the dimensionless conductance as below): calcium current I_{Ca} , potassium current I_K , generic slow current I_S , and background current $I_{K(ATP)}$ through the ATP-dependent potassium channel. The last term, where the prime in the summation stands for the restriction that j be a nearest neighbor of i, represents the electrical coupling via gap junctions of (dimensionless) conductance g_c between nearest neighboring cells in an islet.

The first three types of current depend on the membrane potential V_i and on the gating variables N_i or S_i , and are given by

$$\begin{split} I_{iCa} &= g_{Ca} M_{\infty} (V_i - V_{Ca}) \\ I_{iK} &= g_K N_i (V_i - V_K) \\ I_{iS} &= g_S S_i (V_i - V_K) \end{split}$$

with constant (dimensionless) conductances g_{Ca} , g_K , and g_S and the reversal potentials V_{Ca} and V_K for Ca^{2+} and K^+ ions, respectively. The gating variables N_i and S_i are governed by the equations

$$\tau \frac{dN_i}{dt} = \lambda (N_{\infty} - N_i) \tag{2}$$

$$\tau_S \frac{dS_i}{dt} = S_\infty - S_i \tag{3}$$

where the activation values N_{∞} and S_{∞} as well as M_{∞} in general depend on the membrane potential. Thus the calcium current has been assumed to respond instantaneously to a change in the membrane potential via the voltage-dependent activation M_{∞} while the potassium current is governed by the gating variable N_i via Eq. (2). These two currents are responsible for generating the action potential during the active phase of bursting. The generic slow current, which may corresponds to the inhibitory potassium current, is gated by the slow variable S_i via Eq. (3) and responsible for switching between the active and silent phases. We use Boltzmann-type expressions for the equilibrium values M_{∞} , N_{∞} , and S_{∞} of the voltage-dependent activation:

$$X_{\infty} = \frac{1}{1 + \exp\left[(V_X - V_i)/\theta_X\right]}$$

with appropriate constants V_X and θ_X , where X denotes M, N, or S.

On the other hand, $I_{K(ATP)}$, which is responsible for setting the plateau fraction, i.e., the ratio of the active

phase duration to the burst period, satisfies

$$I_{iK(ATP)} = g_{K(ATP)}p_i(V_i - V_K),$$

where the conductance $g_{K(ATP)}$ depends on the glucose concentration and p_i is the opening probability of the K(ATP) channel. In general, an increase in the intracellular glucose concentration due to injection of glucose raises the ATP/ADP ratio after being metabolized; this in turn reduces $g_{K(ATP)}$ and induces depolarization of the membrane potential [20]. Considering the time delay in this process, i.e., between the increase of the intracellular glucose concentration G_{in} and the reduction of the conductance $g_{K(ATP)}$, we assume that $g_{K(ATP)}$ at time t is a decreasing function of G_{in} at time $t-t_d$, where the delay time t_d measures the time needed for metabolizing glucose transported into the β cell. To be specific, we take the Hill equation [21]:

$$g_{K(ATP)}(t) = \frac{g_1 - g_2}{[G_{in}(t - t_d)/G_K]^b f_D(J_i) + 1} + g_2,$$

with the appropriate Hill coefficient b [22]. Note the sharp contrast to the constant $g_{K(ATP)}$ in existing studies on the bursting mechanism [1] or to the case that $g_{K(ATP)}$ is taken to be a function of ATP [6]. Here g_1 and g_2 represent the maximum and minimum values of $g_{K(ATP)}$, respectively, whereas G_K is the value of G_{in} for which $g_{K(ATP)}$ reduces to its medium value $(g_1 + g_2)/2$. The function $f_D(J_i)$ describes possible direct effects of insulin on $g_{K(ATP)}$. Namely, while we have simply $f_D = 1$ in the absence of such direct effects, in their presence we take $f_D(J_i)$ to be a decreasing function $f(J_i)$ of the inhibition variable J_i :

$$f(J_i) = 1 - \frac{1 + J_0^m}{1 + (J_0/J_i)^m}$$

with the inhibition variable governed by

$$\tau_J \frac{dJ_i}{dt} = J_\infty - J_i. \tag{4}$$

The (slow) relaxation time τ_J and the constant J_0 depend on the pathway, whereas J_{∞} is taken to be an increasing function of the insulin concentration H_i :

$$J_{\infty} = \frac{H_i}{H_J + H_i}$$

with appropriate constant H_J (see Table I).

Although there are many sources of noise in a biological system, we here restrict ourselves only to the stochastic opening and closing of ion channels, in particular those of K(ATP) channels, and assume that the dynamics of the opening probability p_i in the *i*th cell is described by the stochastic differential equation [1, 23]

$$\frac{dp_i}{dt} = \frac{\gamma_1}{\tau_p} (1 - p_i) - \frac{\gamma_2}{\tau_p} p_i + \eta_i(t)$$
 (5)

with appropriate constants γ_1 and γ_2 and relaxation time τ_p , where $\eta_i(t)$ is the Gaussian white noise with zero mean and variance

$$\langle \eta_i(t)\eta_j(t')\rangle = \frac{\gamma_1(1-p_i)+\gamma_2p_i}{\tau_p n_{K(ATP)}}\delta_{ij}\delta(t-t').$$

Note that as the number $n_{K(ATP)}$ of K(ATP) channels per cell grows large $(n_{K(ATP)} \to \infty)$, we have $\eta_i \to 0$ and obtain the usual deterministic equation for p_i .

Modifying the model for a perifusion system [2], we now obtain equations describing glucose regulation. In view of the perifusion, we consider conservation of insulin and write the equation for the insulin concentration H_i around the ith cell in the form:

$$\frac{dH_i}{dt} = \frac{R_s}{\Omega} - k_0(H_i - H_0),$$
 (6)

where R_s is the rate of insulin secretion, k_0 is the flow rate for each cell, and H_0 is the background concentration. The effective volume Ω of a β cell is given by the volume of the islet divided by the number of β cells in it. Experimental observations indicate that insulin secretion from β cells increases with the intracellular calcium concentration and with the amount of insulin stored for rapid secretion [24]. In particular, the secretion rate is known to depend on the Ca²⁺ concentration quartically. We thus assume a quartic function of the concentration C_i of the free intracellular Ca²⁺:

$$R_s = R_0 \tilde{C}_i^4 \frac{S_R}{S_{max}},$$

where R_0 sets the scale, \tilde{C}_i is the rescaled (dimensionless) calcium concentration defined to be $\tilde{C}_i \equiv (C_i - C_b)/C_0$ with the background value C_b and the appropriate scale C_0 , and S_R is the storage amount of insulin preparing rapid secretion in the readily releasable vesicle pool (RRVP) of the β -cell, with the maximum value S_{max} . The latter changes according to the relation

$$\frac{dS_R}{dt} = R_r - R_s,\tag{7}$$

where the refilling rate R_r of insulin granules to the RRVP is proportional to the remaining amount:

$$R_r = a_r (S_{max} - S_R)$$

with appropriate constant a_r . The Ca²⁺ concentration C_i in the *i*th cell is governed by the rate equation:

$$\frac{dC_i}{dt} = -\epsilon(\alpha I_{iCa} - M_S - M_P - M_N + M_{leak}), \quad (8)$$

where ϵ is the fraction of free Ca²⁺ ions in the cytosolic compartments, α is the proportionality constant between the current flow and the concentration reduction rate. Clearance terms, M_S , M_P , and M_N are operated by sarco-endoplasmic reticulum Ca²⁺-ATPase (SERCA)

pumps, plasma-membrane Ca²⁺ ATPase, and plasma-membrane Na/Ca²⁺ exchangers, respectively [25], and given by

$$M_{S} = M_{S}^{max} \frac{1}{1 + (K_{S}/C_{i})^{2}}$$

$$M_{P} = M_{P}^{max} \frac{1}{1 + (K_{P}/C_{i})} \frac{[H^{+}]}{[H^{+}] + [K_{a}]}$$

$$M_{N} = k_{N}C_{i}$$

with appropriate constants M_S^{max} , M_P^{max} , k_N , K_S , and K_P as well as [H⁺] and [K_a] adjusted to pH= 7.40 and pK_a = 7.86 for activation of pumping by protons. The last term M_{leak} represents the flux out of the endoplasmic reticulum (ER) and is described by

$$M_{leak} = \xi_{er}(C_i^{er} - C_i),$$

where ξ_{er} is the rate of Ca²⁺ release from the ER and C_i^{er} denotes the Ca²⁺ concentration in the ER. The latter is governed by the rate equation

$$\frac{dC_i^{er}}{dt} = \epsilon_{er} \frac{\nu_{cyt}}{\nu_{er}} (M_S - M_{leak}), \tag{9}$$

where ϵ_{er} is the fraction of free Ca²⁺ ions in the ER, ν_{cyt} and ν_{er} are the volumes of the cytosolic and of the ER compartments, respectively [26].

Finally, we consider the rate equation for the glucose concentration G_{in} in the *i*th cell:

$$\frac{dG_{in}}{dt} = r_1 + r_2 - r_m, (10)$$

where $r_{1,2}$ are the rates of glucose uptake through GLUT-1 and GLUT-2 transporters, respectively, and r_m is the rate of glucose metabolism [22]. Here the difference between the glucose concentration in blood plasma and the injected glucose concentration G_0 has been disregarded for convenience, with the assumption that the former saturates rather quickly to the latter. We thus take the rates through GLUT-1 and GLUT-2 transporters to be simple increasing functions of the extracellular glucose concentration G_0 and choose an increasing function of the intracellular glucose concentration G_{in} for the metabolism rate:

$$r_{1} = \frac{c_{1}(G_{0} - G_{in})}{(1 + G_{0}/K_{1})(K_{1} + G_{in})} \frac{H_{i}}{H_{r} + H_{i}}$$

$$r_{2} = \frac{c_{2}(G_{0} - G_{in})f_{I}(J_{i})}{(1 + G_{0}/K_{2})(K_{2} + G_{in})}$$

$$r_{m} = \frac{c_{m}G_{in}}{K_{m} + G_{in}}.$$

The constants K_1 , c_2 , K_2 , c_m , and K_m are determined from experimental results, whereas the constants c_1 and H_r are chosen to yield appropriate intracellular glucose concentrations in simulations. Note that r_1 has been

taken to be an increasing function of H_i , which models the recruitment of GLUT-1 transporters by insulin and results in positive feedback. On the other hand, to take into account indirect effects of insulin on $g_{K(ATP)}$, r_2 has been chosen to depend on the function f_I of the inhibition variable J_i , which accounts for inhibitory effects of insulin on its own release via the inhibition of GLUT-2; this assumption is made as a possible interpretation of the recent observation of the inhibiting role of insulin [10, 11, 12]. Accordingly, r_2 tends to decrease as the insulin concentration is raised. Note in particular that unlike the form in the existing study [2], r_2 in the above form always remains positive, thus avoiding the unphysiological case of negative values. To complete the model, we assume the two-step process in a β cell, which includes manifestation of the inhibitory signal from insulin receptors in the presence of insulin (first step) and the inhibition of GLUT-2 transporters by that signal (second step); this provides the slow negative feedback process, which may again be described by the decreasing function $f(J_i)$ (for the direct effects before) with appropriate values of the parameters.

The whole feedback model is thus described by the ten coupled equations [Eqs. (1) - (10)] for the ten variables, together with the equations giving dependence on those variables. It should be stressed that the bursting mechanism and the glucose regulation are incorporated, with the glucose-dependent conductance $g_{K(ATP)}$ and the calcium-dependent rate of insulin secretion as well as the insulin-dependent rate of glucose uptake playing the role of a bridge between the two. The appropriate values of the parameters used in the whole model are given in Table I.

III. SIMULATION RESULTS

We integrate numerically the ten equations, with the parameter values given by Table I, by means of the fifth-order Runge-Kutta method. To keep track of the state of every single channel is computationally expensive; it is sufficient to monitor only the total number of open and closed channels per cell. The time step Δt in the numerical integration should be chosen sufficiently small so that the probability of two channel openings (or closings) during one time step is negligible, say, less than 0.1 [1]:

$$\Delta t < 0.1 \min \left[\frac{\tau_p}{\gamma_1 n_{K(ATP)}}, \frac{\tau_p}{\gamma_2 n_{K(ATP)}} \right].$$

In order to control efficiently the error in the integration, we here use the adaptive step-size algorithm with $\Delta t_{max} = 1$ ms, keeping the above restriction on Δt . This controls automatically the maximum of the next step size to achieve the predetermined accuracy.

TABLE I: Parameter values in the whole feedback model: The parameters without asterisks or single asterisks are fixed or adjustable ones, determined from experiments, respectively. Those with double asterisks are free parameters, not determined from experiments.

$g_{Ca} = 3.6$	$V_{Ca} = 25 \mathrm{mV}$	$V_M = -20 \mathrm{mV}$	$\theta_M = 12 \mathrm{mV}$	$\tau = 20\mathrm{ms}$
$g_K = 10$	$V_K = -75 \mathrm{mV}$	$V_N = -17 \mathrm{mV}$	$\theta_N = 5.6 \mathrm{mV}$	$\lambda = 0.8$
$^*g_S = 4$	$^*\gamma_1=1$	$^*V_S = -22 \mathrm{mV}$	$^*\theta_S = 8.0 \mathrm{mV}$	$^*\tau_S = 60\mathrm{s}$
$\alpha = 1.3 \times 10^{-6} \mathrm{M/Vs}$	$^*\gamma_2=1$	$^*\tau_p = 0.30\mathrm{s}$	$*n_{K(ATP)} = 10^3$	$\Omega = 1.5 \times 10^{-15} \mathrm{m}^3$
$\epsilon = 0.01$	$\epsilon_{er} = 0.01$	$^*\xi_{er} = 1.0 \times 10^{-4}$	$\nu_{cyt} = 10 \mu \mathrm{m}^3$	$\nu_{er} = 0.4 \mu \mathrm{m}^3$
$g_1 = 3.0$	$g_2 = 0.6$	$G_K = 2.8 \mathrm{mM}$	b = 2.5	$^*S_{max} = 3.6 \times 10^{-17} \mathrm{mol}$
$^*M_S^{max}=246\mu\mathrm{M/s}$	$^*M_P^{max} = 126 \mu\mathrm{M/s}$	$^*k_N = 84 \mathrm{s}^{-1}$	$K_S = 0.27 \mu\mathrm{M}$	$K_P = 0.50 \mu\mathrm{M}$
$c_m = 4.1 \times 10^{-5} \mathrm{M/s}$	$K_m = 7.8 \mathrm{mM}$	$^*R_0 = 1.2 \times 10^{-19} \mathrm{mol/s}$	$^*a_r = 2.0 \times 10^{-3} \mathrm{s}^{-1}$	$[H^+] = 10^{-7.40} \mathrm{M}$
$c_1 = 5.8 \times 10^{-4} \mathrm{M/s}$	$K_1 = 1.4 \mathrm{mM}$	$c_2 = 5.3 \times 10^{-4} \mathrm{M/s}$	$K_2 = 17 \mathrm{mM}$	$[K_a] = 10^{-7.86} \mathrm{M}$
$^{**}k_0 = 1.8 \times 10^{-2} \mathrm{s}^{-1}$	$^{**}H_r = 1.4 \times 10^{-2} \mathrm{mM}$	$^{**}H_J = 1.4 \times 10^{-4} \mathrm{mM}$	$^{**}t_d = 90 \mathrm{s}$	$^{**}m = 4.0$
$^{**}H_0 = 0\mathrm{mM}$	$^*C_0 = 10^{-4} \mathrm{mM}$	$^*C_b = 6.0 \times 10^{-5} \mathrm{mM}$		

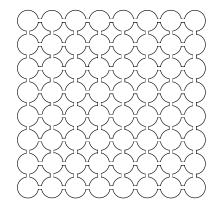


FIG. 2: Two-dimensional square lattice

A. Indirect Pathway

In order to consider indirect effects on $g_{K(ATP)}$ (via negative feedback to GLUT2 by insulin), we take $f_D = 1$ and $f_I = f$ with $J_0 = 0.35$, $\tau_J = 1$ min, and other parameter values in Table I. In this manner we have performed simulations of an islet, which consists of 64 cells arranged on a two-dimensional square lattice shown in Fig. 2, and obtained the membrane potential, Ca²⁺ concentration, and insulin secretion at the coupling conductance $g_c = 0.06$. Figure 3 shows the obtained time dependence of the membrane potential of a single cell randomly chosen among 64 cells for extracellular glucose concentration $G_0 = (a) 3$, (b) 5, (c) 9, and (d) 19 mM. Observed is the bursting behavior of the membrane potential, forming periodic clusters of bursts. Also, the first burst persists longest, while following ones last progressively shorter, similar to experimental observations [24]. The total duration time of such a regularly bursting cluster grows with the glucose concentration, as shown in Fig. 4. Here the criterion for an active phase has been chosen to be the membrane potential higher than $V = -55 \,\mathrm{mV}$; the duration time T represents the total duration of such active

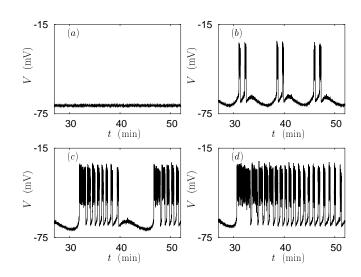


FIG. 3: Bursting behavior of the membrane potential at several extracellular glucose concentrations: $G_0 = (a) 3$, (b) 5, (c) 9, and (d) 19 mM.

phases in a cluster, averaged over clusters. The intracellular glucose concentration G_{in} of a single cell under the same conditions as in Fig. 3 is displayed in Fig. 5; note the oscillatory behavior, with the amplitude and the period growing with the extracellular glucose concentration G_0 .

Figures 6, 7, and 8 exhibit the corresponding behaviors of the concentration of cytosolic calcium Ca²⁺, the insulin secretion rate, and the insulin concentration, respectively, under the same glucose concentrations as in Fig. 3. Unlike other quantities obtained for a single cell, the calcium concentration in Fig. 6 has been averaged over all cells. Note that oscillations of the averaged calcium concentration have periods similar to those of repetitive activation of the membrane potential; this reflects that most of the cells in an islet are well synchronized except for slight phase shift [27], as observed in simulations. It is observed that the Ca²⁺ concentra-

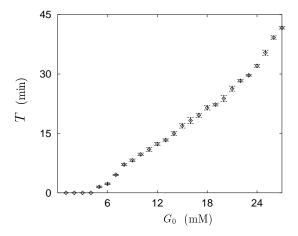


FIG. 4: Duration time T of a bursting cluster depending on the extracellular glucose concentration. Also shown is the typical error bar, estimated from the standard deviation.

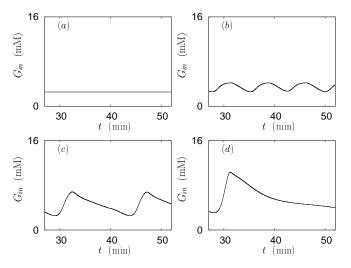


FIG. 5: Behavior of the intracellular glucose concentration at the same extracellular glucose concentrations as in Fig. 3.

tion keeps increasing during the action potential firing and then decreases slowly after the firing stops. The insulin release occurs in the form of sharp peaks as soon as the Ca²⁺ concentration increases in the cell (see Fig. 7). After the rest period of about 5 min, the secretion rate of insulin reaches the maximum value and then diminishes at the next burst, as the amount of insulin stored for rapid secretion reduces. The overall behavior of the insulin secretion rate agrees well with the experimental results [16].

When the flow rate is sufficiently high, on the other hand, the intracellular glucose concentration tends to saturate rather than to oscillate. To probe bursting patterns in this stationary state, we set the insulin flow rate $k_0 = 9.2 \,\mathrm{s}^{-1}$, maintaining previous values of other parameters. Such a low insulin feedback condition may correspond to the experiment performed to obtain stationary values of the intracellular glucose concentration

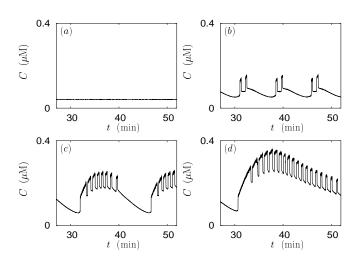


FIG. 6: Behavior of the calcium concentration at the same extracellular glucose concentrations as in Fig. 3.

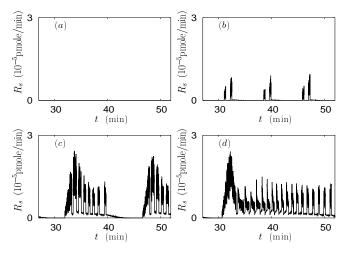


FIG. 7: Behavior of the insulin secretion rate at the same extracellular glucose concentrations as in Fig. 3.

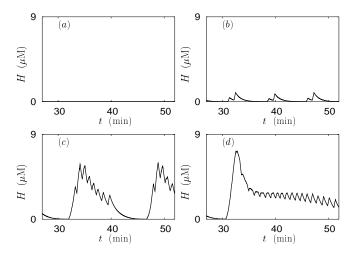


FIG. 8: Behavior of the insulin concentration around a cell at the same extracellular glucose concentrations as in Fig. 3.

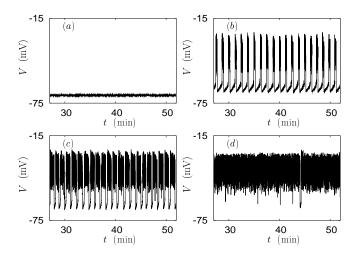


FIG. 9: Behavior of the membrane potential in the stationary state of the intracellular glucose concentration. The extracellular glucose concentrations are the same as those in Fig. 3.

at several extracellular glucose concentrations [28]. The resulting behavior of the membrane potential is displayed in Fig. 9, depending on the extracellular glucose concentration. The duration time of each burst is also observed to grow with the extracellular glucose concentration; on the other hand, oscillatory patterns with periods of about 10 min do not emerge. It is thus concluded that the emergence of slow oscillations of the membrane potential depends upon the flow rate of the circulating extracellular fluid, i.e., blood. The lack of slow oscillations led by minimal insulin feedback matches well with the experimental results of IRS-1 (insulin receptor substrate 1) knockout mice [29].

Figure 10 shows how the intracellular glucose concentration G_{in} depends on the extracellular glucose concentration G_0 . Shown together with the experimental results [28] are the saturated values of G_{in} obtained from simulations with G_0 varied from 1 mM to 26 mM. It is pleasing to see the excellent agreement between the experimental data and the simulation data. In particular, the dependence of G_{in} on G_0 yields reasonably good linear fitting, with the slope 0.81.

We also investigate the case that the concentration of injected glucose is varied (with $k_0 = 9.2\,\mathrm{s}^{-1}$), to probe time-dependent responses to sudden changes of the extracellular glucose concentration. For simplicity, the extracellular glucose concentration G_0 is given as a step function of time, with possible time delay disregarded. Figure 11 shows the behaviors of the membrane potential, calcium concentration, and intracellular glucose concentration as glucose is injected with the concentration $G_0 = 5, 9, 12, 19$, and 11 mM at time t = 20, 30, 40, 50, and 60 min, respectively. They are observed to display behaviors in close response to the change of the glucose concentration.

To confirm the assertion that insulin inhibits bursting of membrane potentials, we examine the behavior of

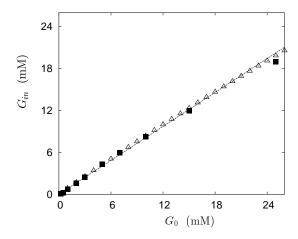


FIG. 10: Saturated intracellular glucose concentration depending on the extracellular glucose concentration. Squares represent the data obtained from experiment [28], whereas triangles are the data obtained via simulations and fitted by the dashed line.

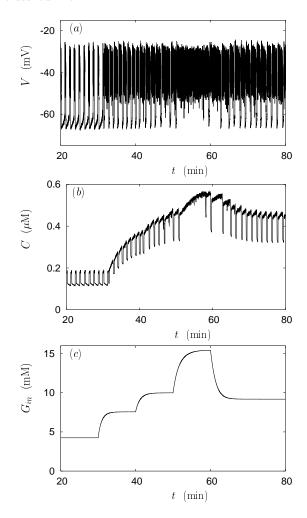


FIG. 11: Behavior of (a) the membrane potential, (b) calcium concentration, and (c) intracellular glucose concentration when glucose is injected as a step function of time: $G_0 = 5, 9, 12, 19,$ and $11 \,\mathrm{mM}$ at time t = 20, 30, 40, 50, and $60 \,\mathrm{min}$, respectively.

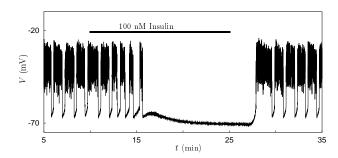


FIG. 12: Behavior of the membrane potential under $G_0 = 10 \,\mathrm{mM}$, when insulin corresponding to $H_0 = 100 \,\mathrm{nM}$ is injected during the time $t = 10 \,\mathrm{min}$ to 25 min. The experimental values $c_1 = 2.9 \times 10^{-4} \,\mathrm{M/s}$ and $k_0 = 9.2 \,\mathrm{s}^{-1}$ are used.

the membrane potential in response to the injection of insulin. Assuming that the basal plasma insulin concentration changes rapidly, we represent the insulin stimulus by an increase in the basal insulin concentration H_0 . Figure 12 shows the membrane potential in response to the insulin injection corresponding to $H_0 = 100 \,\mathrm{nM}$ during the time $t = 10 \,\mathrm{min}$ to 25 min under the glucose concentration $G_0 = 10 \,\mathrm{mM}$. (Note that the parameter values for normal β cells are used here for comparison, different from those for hamster insulinoma tumor cells in Table I. See Ref. 2 and references therein.) It is observed that the bursting action potential disappears after about 6 min from the beginning of the insulin injection and appears again in about 3 min after the cease of injection. This is consistent with the experimental result [11].

Finally, for a more realistic model, we allow variations in some parameters among the cells, and examine the effects of such inhomogeneity. We have considered variations in the coupling conductance g_c , choosing randomly its value in the range $0.05 < g_c < 0.07$. The resulting behavior of the membrane potential is found to be qualitatively the same as that for constant g_c .

B. Direct Pathway

There are reports that insulin itself inhibits insulin secretion in β -cells although the exact mechanism is not clear [10, 11]. Specifically, experimental data seem to suggest that the negative feedback of insulin acts directly on ATP-dependent K⁺ channels, suppressing action potential firing [11]. To examine this direct pathway, we set $f_I = 1$, $f_D = f$ with $J_0 = 0.15$, and $H_J = 1.0 \times 10^{-3}$ mM, varying the relaxation time τ_J . Figure 13 exhibits the resulting oscillatory behaviors of the membrane potential for $\tau_J = 3 \,\mathrm{min}$, depending on the extracellular glucose concentration G_0 . The average duration of bursting is rather insensitive to G_0 , approximately 16 seconds for a wide range of G_0 . Compared with the results of the indirect pathway, the duration time is short and the doseresponse to glucose is also weak in the direct pathway. On the other hand, the average period between bursts

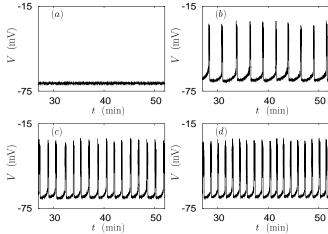


FIG. 13: Bursting behavior of the membrane potential via the direct pathway mechanism with $\tau_J = 3 \,\text{min}$, at the same extracellular glucose concentrations as in Fig. 3.

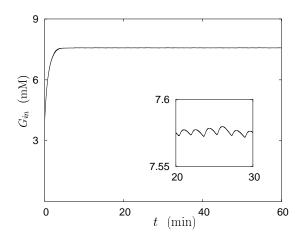


FIG. 14: Behavior of the intracellular glucose concentration at $G_0 = 9\,\mathrm{mM}$ and $\tau_J = 3\,\mathrm{min}$ in the presence of the direct pathway. The inset shows an enlarged view.

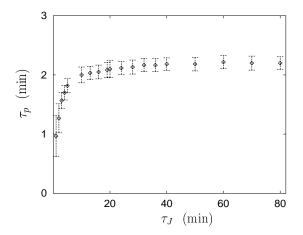


FIG. 15: Bursting period τ_p depending on τ_J at $G_0 = 10$ mM.

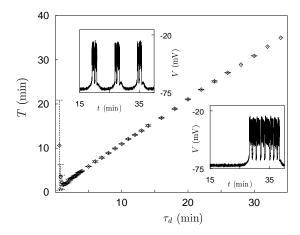


FIG. 16: Duration time T of a bursting cluster versus the time delay τ_d , with typical error bars estimated from the standard deviation. The insets exhibit the bursting patterns at $\tau_d = 1.2 \, \text{min}$ (upper left) and 9.0 min (lower right).

decreases with G_0 , which results in more insulin secretion at higher glucose levels. Unlike the indirect pathway, however, the clustering of bursts does not emerge in the direct pathway. Since the secreted insulin affects the ATP-dependent K⁺ channels directly and rapidly, there is not enough time to activate the mechanism for bursting clusters. The behavior of the intracellular glucose concentration G_{in} is shown in Fig. 14, for the extracellular glucose concentration $G_0 = 9 \,\mathrm{mM}$ and the relaxation time $\tau_J = 3 \,\mathrm{min}$. No prominent slow oscillations are observed (compare with Fig. 5). Figure 15 shows how the average bursting period τ_p depends on the relaxation time τ_J : Observed is a rapid increase of τ_p , followed by saturation to $2 \,\mathrm{min}$, as τ_J is raised.

To obtain bursting clusters via the direct pathway, one needs another formal time delay and accordingly, an additional process in insulin inhibition. As an attempt, we consider delay τ_d in the direct activation factor, i.e., $f_D(t-\tau_d)$ on the conductance of the ATP-dependent K⁺ channel, and show the results in Fig. 16. It is observed that the total duration time T of a bursting cluster grows almost linearly with the delay τ_d . When τ_d is small, we have $T \approx 0$ and there is no regular bursting patterns. When τ_d is raised to 1.2 min, clustering patterns of bursts begin to appear; the corresponding behavior of the membrane potential is displayed in the upper inset panel. As τ_d is increased further, the bursting duration time keeps growing and the bursting pattern at $\tau_d = 9.0 \,\mathrm{min}$ becomes similar to that via the indirect pathway shown in Fig. 3(c). Nevertheless it is rather difficult to recognize this direct pathway as an appropriate mechanism for bursting clusters since such large delay is apparently not consistent with experimental results [11].

IV. DISCUSSION

Our model has employed the glucose-insulin feedback for slow dynamics, responsible for slow oscillations in the system. Unlike the fast oscillations of bursting action potentials, however, the origin of the slow oscillations with periods about 5 to 10 minutes is not understood well and still in controversy. One of the main difficulties lies in identifying the primary oscillation among many secondary ones. Indeed in a pancreas, many chemical ingredients, e.g., Ca²⁺, O₂, ATP, glucose, and insulin, show oscillatory behaviors with periods of several minutes [8, 17, 18, 30]. Recently, it has been reported that glycolysis produces spontaneous oscillations in β -cells: Here Ca²⁺ acts as a mediator, transducing oscillatory metabolism into oscillatory secretion [30]. After completion of the first version of the manuscript, we became aware of the recent work, which addresses the multiple bursting mode in pancreatic islets, with glycolysis adopted as the slow dynamics [31]. In the present study, on the other hand, we have proposed another possible mechanism for slow oscillations, namely, the negative feedback through the product (i.e., insulin) in response to the stimulus (glucose).

In combining the slow dynamics (glucose-insulin feedback) with the fast dynamics (bursting action potentials), we have allowed both dynamics to interact bidirectionally. As a result, there emerge mixed oscillations of both time scales, clusters of bursts, when the two dynamics complement each other via the indirect pathway. This also keeps parallel with the glycolytic model [31], and contrasts with the unidirectionally interacting model where slow dynamics of glycolysis influences fast dynamics one-sidedly [32].

We now discuss the plausibility of our model, in comparison with others. Insulin secretion is synchronized among the β cells in an islet, which has been observed in the calcium imaging experiment [27]. In addition, the pulsatile insulin secretion in a group of islets indicates additional synchronization among islets in a pancreas; otherwise the secreted insulin from many islets, having different phases from each other, would produce essentially a constant level of insulin [33], although there may exist a source of some controversy [34]. Such synchronization among islets may not be explained by the model adopting glycolysis as the slow dynamics, where no communication is available between islets. An expanded model was also proposed, incorporating glucose release in the liver: Secreted insulin causes oscillations of the plasma glucose through the liver, which in turn entrain all islets with its characteristic frequency [33]. This is, however, in disagreement with pulsatile insulin secretion at constant glucose concentration observed in vitro [35] and in vivo [36]. On the other hand, in our feedback model, secreted insulin from an islet affects other islets (via the paracrine interaction) as well as the islet itself (via the autocrine interaction), thus inherently allowing synchronization among islets even at a constant plasma glucose level. Accordingly, our model can explain the sustained insulin oscillations at constant glucose concentration observed in vitro and in vivo [35, 36]. Further, the recent observation that defects in IRS-1, associated with the reduced expression of the SERCA protein, result in lack of slow calcium oscillations and reduction of insulin secretion [29] is also supportive of the insulin feedback mechanism. Note, however, that the molecules for negative feedback are not necessarily limited to insulin and other products, e.g., GABA, can also provide the feedback pathway [37].

In constructing a model for such complex biological systems, one may not expect perfection and impeccableness, and there are certainly some drawbacks in our feedback model as well. First, our model predicts that both the amplitude and the period of the insulin oscillation increase with the plasma glucose concentration, which apparently does not agree with the report that mainly the amplitude tends to change [18, 36]. Secondly, slow oscillations are still observed in experiment on single β cells, even when the flow rate of the local perfusion is too fast for secreted insulin to reach the insulin receptor and to operate the feedback mechanism [38]. In this case the other model adopting glycolysis as slow dynamics is still applicable since the glycolytic oscillation may still occur without the feedback through a product like insulin [31].

In summary, we have incorporated the macroscopic description of glucose regulation and the microscopic mechanism of bursting behavior of β cells, to establish a model that displays inherently the observed oscillations of the membrane potential, cytosolic calcium concentration, and insulin secretion in pancreatic islets. In view of the experimental observations, the ATP-dependent conductance $g_{K(ATP)}$ of the K⁺ channel has been taken as a decreasing function of the glucose concentration whereas the insulin secretion rate has been given by an increasing function of the intracellular Ca²⁺ concentration. Considered here are two possible insulin induced inhibitory

pathways, affecting the conductance $g_{K(ATP)}$. Further, β cells in an islet have been considered to be coupled electrically with their nearest neighbors with appropriate coupling conductance. By means of extensive numerical simulations, we have obtained bursting electrical activities, calcium concentration, and insulin secretion, which are consistent with those observed experimentally. In addition to the stationary-state behaviors, we have also examined how the intracellular glucose concentration changes with the extracellular glucose concentration, and explored the behaviors of the system as the injected glucose concentration is varied successively. Then the inhibition effects of insulin on the bursting action potential have been probed. Finally, we have compared two possible pathways to the feedback modulations of $g_{K(ATP)}$ by secreted insulin. With appropriate parameters, the indirect feedback mechanism has been found to generate bursting clusters of action potentials. On the other hand, the direct mechanism in general generates only single bursts; it requires, e.g., rather unrealistically long delay time to produce bursting clusters.

Our model may thus serve as a useful framework to predict insulin secretion modulated by the molecular mechanism in β cells and the corresponding glucose regulation. In particular, our model, based on the glucose-insulin feedback, can be a good complement to the glycolytic model. Extensions toward a more realistic description of the biological situation are left for further study.

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G. de Vries and A. Sherman, J. Theor. Biol. 207, 513 (2000).

^[2] L.W. Maki and J. Keizer, Am. J. Physiol. 268, C780 (1995); Bull. Math. Biol. 57, 569 (1995).

^[3] J. Sturis, K.S. Pololnsky, E. Mosekilde, and E.V. Cauter, Am. J. Physiol. 260, E801 (1991).

 ^[4] G. de Vries, Phys. Rev. E 64, 051914 (2001); N.F.
 Rulkov, Phys. Rev. Lett. 86, 183 (2001); Phys. Rev. E 65, 041922 (2002).

^[5] A.L. Hodgkin and A.F. Huxley, J. Physiol. (London) 117, 500 (1952).

^[6] J. Keizer and G. Magnus, Biophys. J. 56, 229 (1989); P. Smolen and J. Keizer, J. Membrane Biol. 127, 9 (1992).

^[7] P.J. Lefébvre, G. Paolisso, A.J. Scheen, and J.C. Henquin, Diabetologia 30, 443 (1987); R.M. Barbosa, A.M. Silva, A.R. Tomé, J.A. Stamford, R.M. Santos, and L.M. Rosário, J. Physiol. 510, 135 (1998).

^[8] Y.J. Liu, A. Tengholm, E. Grapengiesser, B. Hellman,

and E. Gylfe, J. Physiol. (London) 508, 471 (1998).

^[9] F.M. Ashcroft, P. Proks, P.A. Smith, C. Ammala, K. Bokvist, and P. Rorsman, J. Cell. Biochem. 55 (Suppl.), 54 (1994).

^[10] S.J. Persaud, H. Asare-Anane, and P.M. Jones, FEBS Letters 510, 225 (2002).

^[11] F.A. Khan, P.B. Goforth, M. Zhang, and L.S. Satin, Diabetes 50, 2192 (2001).

^[12] C.S. Nunemaker, M. Zhang, and L.S. Satin, Diabetes 53, 1765 (2004).

^[13] J.V. Sánchez-Andrés, A. Gomis, and M. Valdeolmillos, J. Physiol. (London) 486, 223 (1995); A. Gomis, J.V. Sánchez-Andrés, and M. Valdeolmillos, Pflügers Arch. 432, 510 (1996).

^[14] S. Dryselius, E. Grapengiesser, B. Hellman, and E. Gylfe, Am. J. Physiol. 276, E512 (1999).

^[15] B. Hellman, and E. Gylfe, and E. Grapengiesser, P.E. Lund, and A. Berts, Biochim. Biophys. Acta 1113, 295

- (1992); B.A. Cunningham, J.T. Deeney, C.R. Bliss, B.E. Corkey, and K. Tornheim, Am. J. Physiol. **271**, E702 (1996); P. Bergsten, *ibid.* **274**, E796 (1998); I. Nakagaki, S. Sasaki, S. Hori, and H. Kondo, Pflügers Arch. **440**, 828 (2000).
- [16] P. Gilon, R.M. Shepherd, and J.C. Henquin, J. Biol. Chem. 268, 22265 (1993).
- [17] H.F. Chou and E. Ipp, Diabetes **39**, 112 (1990).
- [18] N. Pørksen, Diabetologia **45**, 3 (2002).
- [19] Here glucose diffusion effects have been considered in J. P. Keener, Bull. Math. Biol. 63, 625 (2001).
- [20] S. Göpel, T. Kanno, S. Barg, J. Galvanovskis, and P. Rorsman, J. Physiol. (London) 521, 717 (1999).
- [21] J.N. Weiss, FASEB J. 11, 835 (1997).
- [22] For simplicity of notation, the subscript i labeling the cell has been omitted.
- [23] L.J. DelFelice and A. Isaac, J. Stat. Phys. **70**, 339 (1992);
 R.F. Fox and Y.N. Lu, Phys. Rev. E **49**, 3421 (1994).
- [24] P. Rorsman and E. Renström, Diabetologia 46, 1029 (2003).
- [25] L. Chen, D.S. Koh, and B. Hille, Diabetes 52, 1723 (2003).
- [26] R. Bertram and A. Sherman, Biophys. J. 87, 3775 (2004).
- [27] R.M. Santos, L. M. Rosario, A. Nadal, J. García-Sancho, B. Soria, and M. Valdeolmillos, Pflügers Arch. 418, 417 (1991).
- [28] R.R. Whitesell, A.C. Powers, D.M. Regen, and N.A.

- Abumrad, Biochem. **30**, 11560 (1991).
- [29] R.N. Kulkarni, M.G. Roper, G. Dahlgren, D.Q. Shih, L.M. Kauri, J.L. Peters, M. Stoffel, and R.T. Kennedy, Diabetes 53, 1517 (2004).
- [30] K. Tornheim, Diabetes 46, 1375 (1997); R.T. Kennedy, L.M. Kauri, G.M. Dahlgren, and S.K. Jung, Diabetes 51, S152 (2002).
- [31] R. Bertram, L. Satin, M. Zhang, P. Smolen, and A. Sherman, Biophys. J. 87, 3074 (2004).
- [32] K. Wierschem and R. Bertram, J. Theor. Biol. 228, 513 (2004).
- [33] M.G. Pedersen, R. Bertram, and A. Sherman, Biophys. J. 89, 107 (2005).
- [34] M. Valdeolmillos, A. Gomis, and J.V. Sánchez-Andrés, J. Physiol. (London) 493, 9 (1996).
- [35] J.I. Stagner, E. Samols, and G.C. Weir, J. Clin. Invest. 65, 939 (1980).
- [36] H.S. Song, S.S. McIntyre, H. Shah, J.D. Veldhuis, P.C. Hayes, and P.C. Butler, J. Clin. Endocrinol. Metab. 85, 4491 (2000).
- [37] M. Braun, A. Wendt, K. Buschard, A. Salehi, S. Sewing, J. Gromada, and P. Rorsman, J. Physiol. 559, 397 (2004).
- [38] E. Grapengiesser, E. Gylfe, and B. Hellman, Biochem. Biophys. Res. Commun. 151, 1299 (1988).